

A Strong Constraint on Radiative Forcing of Well-mixed Greenhouse Gases

Jing Feng

jing.feng@princeton.edu

Princeton University <https://orcid.org/0000-0003-1144-622X>

David Paynter

GFDL

Raymond Menzel

Geophysical Fluid Dynamics Laboratory

Ryan Kramer

<https://orcid.org/0000-0002-9377-0674>

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A Strong Constraint on Radiative Forcing of Well-mixed Greenhouse Gases

Jing Feng^{1*}, David Paynter², Raymond Menzel², Ryan Kramer²

^{1*}Atmospheric and Oceanic Sciences Program, Princeton University,
300 Forrestal Road, Princeton, 08540, New Jersey, United States.

²Geophysical Fluid Dynamics Laboratory, 201 Forrestal Road,
Princeton, 08540, New Jersey, United States.

*Corresponding author(s). E-mail(s): jing.feng@princeton.edu;
Contributing authors: david.paynter@noaa.gov;
raymenzell324@gmail.com; ryan.kramer@noaa.gov;

Abstract

Radiative forcing from well-mixed greenhouse gases (WMGHGs) is a primary driver of Earth's energy imbalance but remains difficult to constrain due to its dependence on atmospheric state and model discrepancies. We present the first global benchmark of longwave (LW) instantaneous radiative forcing (IRF) for major WMGHGs under realistic, all-sky conditions. This benchmark is supported by (1) global line-by-line radiative transfer simulations using 24 years of reanalysis-based atmospheric and cloud conditions, and (2) a robust linear relationship between LW IRF and outgoing longwave radiation, enabling direct observational constraints from satellites. Using this framework, we estimate that the rising WMGHG concentrations have increased LW IRF by $3.71 \pm 0.07 \text{ W/m}^2$ from 1850 to 2024. Applying this benchmark to Earth System Models, we reduce the uncertainty in CO₂ effective radiative forcing by 60%. This observation-informed, physically grounded method offers a more accurate and consistent constraint on greenhouse gas forcing for future climate assessment reports.

Keywords: Radiative Forcing, Greenhouse Gases, Earth System Modeling

27 Main

28 Earth’s radiative energy balance is fundamentally regulated by greenhouse gases
29 (GHGs), which absorb and re-emit infrared radiation, exerting a profound influence on
30 the planet’s climate system. Anthropogenic emissions of well-mixed greenhouse gases
31 (WMGHGs), including CO₂, CH₄, N₂O, chlorofluorocarbons (CFCs) and hydroflouro-
32 carbons (HFCs), perturb this balance by reducing outgoing longwave radiation (OLR),
33 thereby inducing global warming until a new equilibrium is achieved. The magnitude
34 of this perturbation is characterized by instantaneous radiative forcing (IRF), defined
35 as an instantaneous change in net radiative fluxes at top-of-atmosphere (TOA) or
36 tropopause [1–7].

37 Recent decades have brought major advances in our ability to constrain the radia-
38 tive forcing of WMGHGs. Long-term gas concentration records from observational
39 networks such as NOAA Earth System Research Laboratories (ESRL) and NASA
40 Advanced Global Atmospheric Gases Experiment (AGAGE) [8, 9], combined with
41 improvements in molecular spectroscopy and line-by-line radiative transfer modeling
42 [10, 11], have enabled high-fidelity benchmark estimates of IRF for selected atmo-
43 spheric profiles [12–14]. Under carefully sampled conditions, line-by-line calculations
44 have provided high-confidence estimates of IRF for clear-sky (cloud free) conditions.
45 For example, the most recent benchmark for the clear-sky longwave (LW) IRF due
46 to WMGHG increases from 1850 to 2014 is 2.66 to 2.69 Wm^{−2} at TOA, and 3.63 to
47 3.66 Wm^{−2} at tropopause, with sampling errors below 0.015 Wm^{−2} relative to the
48 2014 climatology (Table S1)[14]. The benchmark has not been updated in a decade,
49 despite substantial increases in WMGHG concentrations since 2014 [15].

50 Despite these advances, benchmark estimates of radiative forcing remain limited
51 to fixed clear-sky background conditions and do not capture the broader uncertain-
52 ties associated with evolving surface temperatures, stratospheric states, water vapor,
53 and cloud distributions for CO₂ [16–21] and other WMGHGs [22–27]. Over 90% of

WMGHG IRF arises from the longwave (LW) spectrum [13, 14], where radiative
fluxes are highly sensitive to background atmospheric conditions [28, 29]. line-by-line
radiative transfer calculations using a limited set of observed cloud profiles can esti-
mate all-sky IRF [13, 23], but sampling errors become more uncertain due to the
spatio-temporal variability of clouds [30]. Moreover, when used as inputs to line-by-
line calculations, persistent biases and uncertainties in humidity and cloud fields in
reanalyses [31, 32] and Earth system models (ESMs) [29, 33] introduce additional
uncertainties that are challenging to constrain.

In ESMs, the effective radiative forcing of WMGHGs remains uncertain by
more than 12% [7, 34], with more substantial spread ($\sim 15\%$) in the LW IRF
component [17, 20, 35]. These uncertainties stem from both inter-model differences in
atmospheric state representations [18, 20] and limitations in radiative parameteriza-
tions [14]. Although line-by-line methods applied to various evolving background states
could clarify these inter-model differences, they remain computationally prohibitive.

Satellite measurements of radiance and irradiance show WMGHG forcing signa-
tures [36, 37], but these signals intermingle with variability in surface temperature,
water vapor, clouds, aerosols, and short-lived climate forcers. Isolating the greenhouse
gas component requires radiative transfer simulations based on retrieved or assumed
background conditions [36–40], introducing uncertainties often exceed 10%—compa-
rable to or even greater than the inter-model spread. Although these observations
provide compelling evidence of anthropogenic influence, they have not yet yielded
robust quantitative constraints on IRF.

Together, these limitations underscore a central challenge: despite advances in
spectroscopy, satellite observations, and radiative modeling, the state dependence of
LW IRF remains a leading source of uncertainty in the radiative forcing of greenhouse
gases.

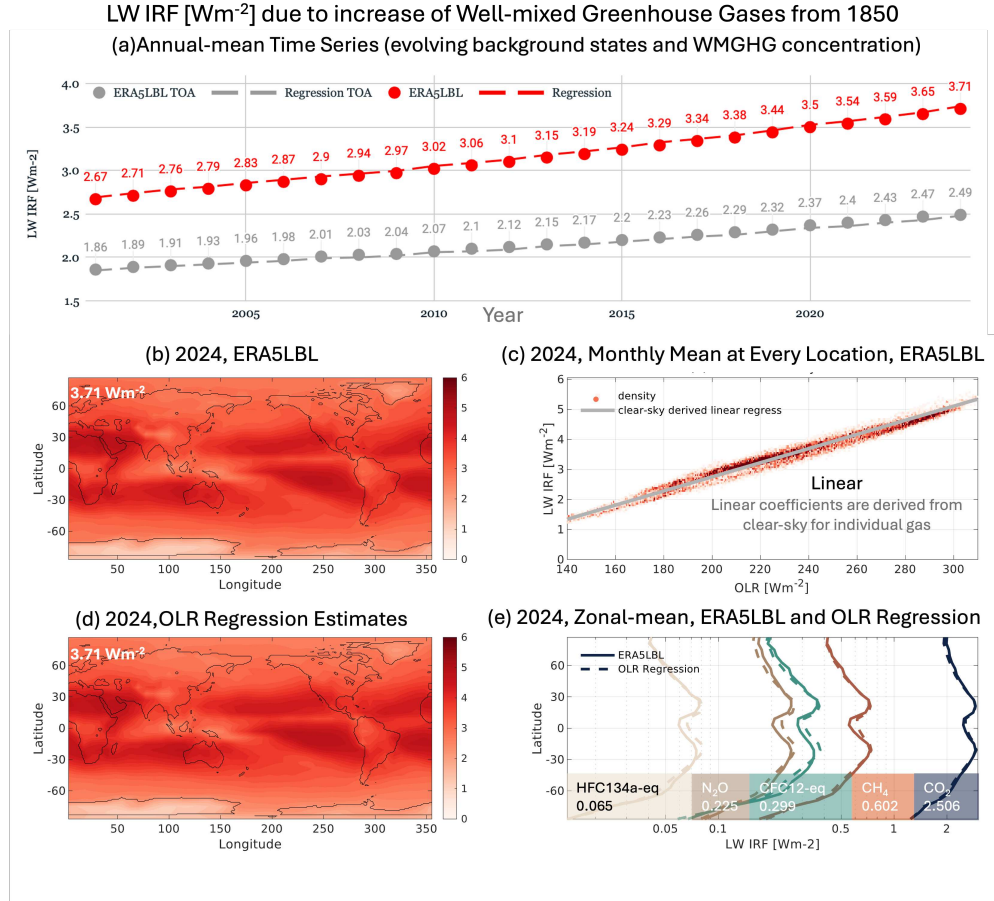


Fig. 1 Global line-by-line calculations reveal a strong linear relationship between LW IRF and OLR. This forms the basis of a regression model (Eq. 1, Table S2 and S3). Panel (a) shows the global annual-mean LW IRF (red: tropopause, gray: TOA) time series from ERA5LBL (solid) and regression estimates based on evolving WMGHG concentrations and OLR from ERA5LBL (dashed). Panels (b–e) evaluate LW IRF at the tropopause from WMGHG increases (1850–2024) under 2024 ERA5LBL background conditions: (b) annual-mean LW IRF from ERA5LBL; (c) joint histogram of monthly LW IRF (y-axis) and OLR (x-axis) compared with regression estimate from clear-sky relationships (gray); (d) LW IRF reconstructed from spatial OLR of ERA5LBL using the regression model; (e) zonal-mean LW IRF from ERA5LBL (solid) and regression applied to zonal-mean OLR of ERA5LBL (dashed) for major WMGHGs.

80 Global-Scale Line-by-Line Benchmark of Longwave

81 Radiative Forcing

82 Enabled by the GPU-optimized line-by-line radiative transfer code (GRTcode) devel-
 83 oped at NOAA’s Geophysical Fluid Dynamics Laboratory, this study presents the

first global-scale, decadal line-by-line calculations of longwave radiative forcing for 84
CO₂, CH₄, N₂O, CFCs, and HFCs. The code is benchmarked by Pincus et al. (2020) 85
[14]. These calculations use monthly-mean surface temperature, air temperature, 86
water vapor, ozone, and cloud from the ERA5 reanalysis dataset [41]. Annual-mean 87
WMGHG concentrations follow the Coupled Model Intercomparison Project Phase 7 88
(CMIP7) input datasets from 2001 to 2022 [9], with CO₂, CH₄, and N₂O extended to 89
2024 using NOAA ESRL trends [42, 43], as detailed in the *Methods*. Following Mein- 90
shausen et al. (2017)[8], radiative effects of ozone-depleting substances are combined 91
as CFC-12 equivalence (CFC12eq), and other fluorinated gases as HFC-134a equiva- 92
lence (HFC134a-eq). Results generated following this experiment setup is referred to 93
as ERA5LBL. 94

The ERA5LBL dataset provides gridded, monthly-mean outgoing longwave radi- 95
ation (OLR) from 2001 to 2024 for a ‘Control’ experiment. To isolate the LW IRF 96
of WMGHG x , relative to pre-industrial (PI, 1850) concentrations, sensitivity experi- 97
ments (x PI) are conducted that are identical to the Control experiment, except that 98
gas x is held fixed at its PI concentration. Additional experiments are performed for 99
doubled and quadrupled PI CO₂ concentrations. Details are provided in the Sup- 100
plementary Materials (*Methods*). For each gas x , the LW IRF at vertical level i 101
is computed as the difference in net downward LW flux between the Control and 102
perturbation (x PI) experiments: 103

$$F_{i,x} = (R_{\downarrow,i,xPI} - R_{\uparrow,i,xPI}) - (R_{\downarrow,i,Control} - R_{\uparrow,i,Control}), \quad (1)$$

where $R_{\uparrow,i}$ and $R_{\downarrow,i}$ denote the upwelling and downwelling LW radiative fluxes 104
at level i . While the IRF at the TOA is often used to decompose the TOA energy 105
budget, the tropopause value is more relevant for surface temperature responses and 106
adjusted forcing [1, 2, 5, 6, 44, 45]. In this paper, LW IRF at TOA is referred to as 107

108 IRF_{TOA}, and ‘LW IRF’ specifically refers to F_i at the **200 hPa** level [4], which serves
109 as a consistent tropopause level for benchmark evaluations.

110 From 2001 to 2024, ERA5LBL simulations show an increase in LW IRF from 2.67
111 to 3.71 Wm⁻², and at TOA from 1.86 to 2.49 Wm⁻² (Fig.1a). The increase in CO₂
112 from 370.8 to 423.0 parts per million by volume (ppmv) accounts for 0.853 Wm⁻² of
113 the increase, while CH₄, N₂O, HFC134a-eq, and CFC12-eq contribute 0.080, 0.073,
114 0.051, and -0.008 Wm⁻², respectively.

115 Compared to the recent benchmark by Pincus et al. (2020) [14], which esti-
116 mated LW IRF_{TOA} from 100 ERA5 profiles in 2014, our global clear-sky LW IRF
117 (3.64 Wm⁻²) and LW IRF_{TOA} (2.67 Wm⁻²) show excellent agreement (Table S1).
118 Pincus et al. did not provide a benchmark for all-sky conditions but estimated LW
119 IRF_{TOA} between 1.98 and 2.07 Wm⁻², based on the all-sky to clear-sky IRF ratio
120 from three CMIP6 models [39, 46, 47]. Our ERA5LBL-derived all-sky LW IRF_{TOA}
121 for 2014 is 2.17 Wm⁻², exceeding the upper bound of this range, suggesting weaker
122 LW cloud effects in ERA5 than those CMIP6 models.

123 Alternative estimates of all-sky LW IRF can be obtained from simplified expres-
124 sions [7, 12, 13] (Table S3), using conversion factors from line-by-line calculations
125 with sampled cloud profiles [48, 49] (Table S4). Combining IPCC AR6[7] for CFC-
126 s/HFCs with Etminan et al. (2016)[13] for CO₂, CH₄, and N₂O gives 3.14 Wm⁻² for
127 2014; using Byrne and Goldblatt (2014)[12] yields 2.99 Wm⁻². Both are lower than
128 the ERA5LBL benchmark (3.19 Wm⁻²), likely due to uncertainties in the conversion
129 factors, simplified expressions, or cloud effects.

130 While ERA5 likely provides more realistic atmospheric and cloud conditions by
131 assimilating all-sky satellite observations [41], potentially improving persistent biases
132 in earlier reanalysis datasets [32], simulating cloud radiative effects in line-by-line
133 models remains sensitive to assumptions about subcolumn cloud inhomogeneity [30,

50, 51], vertical cloud overlap [52, 53], and hydrometeor particle sizes [54–56]. These sensitivities lead to biases and uncertainties that are challenging to quantify.

Constraining Present-day Radiative Forcing Using Robust Linear Dependence on OLR

Despite potential uncertainties in all-sky forcing estimates, the LW IRF of different WMGHGs exhibits highly consistent spatial patterns (Fig.1b; see also Fig.S1), largely independent of cloud conditions. This consistency aligns with the inhomogeneous forcing structure previously identified for CO₂[18]. We find that the spatial variability in IRF closely follows local outgoing longwave radiation (OLR) under both clear-sky and all-sky conditions (Figs.S1-S2). This strong correspondence suggests that temperature, humidity, and clouds influence OLR[28, 29] and LW IRF in similar ways: IRF tends to be higher in regions with warmer surface temperatures [28] and lower where relative humidity or cloud amount is higher [29].

Motivated by this relationship, we construct a simple regression model for each WMGHG using ERA5LBL results under clear-sky conditions:

$$F = a(R - b) - r, \quad (2)$$

where F is the LW IRF, R is the clear-sky OLR under present-day WMGHG concentrations (“Control”; see *Methods*), and a and b are regression coefficients representing the slope and intercept (Fig.S3). The coefficients are derived under the constraint of zero global-mean bias in clear-sky IRF for the year 2010 ($r = 0$) and are summarized in Table S2.

This regression approach accurately reproduces the clear-sky LW IRF spatial pattern for all gases (Figures S2, for 24-year-average) and achieves global-mean accuracy comparable with five benchmark line-by-line models (Table S1)[14]. IRF induced by

individual WMGHGs is additive, with a small global-mean residual of -0.017 Wm^{-2} ($\sim 0.5\%$) in the year 2010.

Remarkably, this linear relationship remains the same under all-sky conditions. By substituting all-sky OLR into Equation 2, the clear-sky-derived regression model accurately predicts the all-sky LW IRF–OLR relationship (gray line in Fig.1c; $R^2 = 0.99$), the spatial distribution of LW IRF (Fig.1d), and the zonal-mean profiles for each gas (Fig.1e; dashed vs. solid curves). The global-mean bias (r in Eq.2) under all-sky conditions remains within 0.01 Wm^{-2} (Table S2). Further evaluations for individual gases and IRF_{TOA} results are provided in the *Methods* (Table S5, Figs. S3–S4).

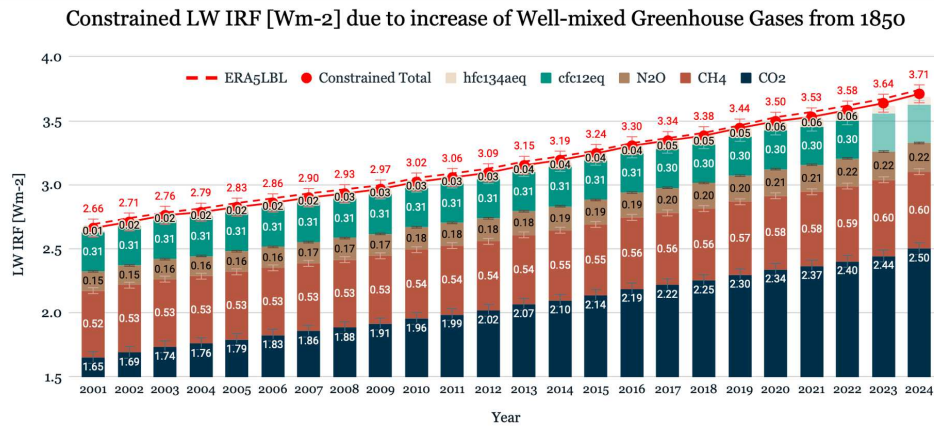


Fig. 2 Constrained LW IRF due to the increase in major WMGHG concentrations relative to 1850, based on clear-sky-derived linear coefficients from ERA5LBL, global-mean timeseries of OLR from CERES EBAF v4.2, and global-mean timeseries of WMGHG concentrations. LW IRF from all-sky ERA5LBL is shown as dashed red curve.

Applying this regression model to time-varying WMGHG concentrations and global-mean OLR from ERA5LBL yields excellent agreement with directly computed LW IRFs (dashed lines versus markers in Fig.1a). The 5–95% confidence range for the global annual-mean LW IRF is conservatively estimated at 0.064 Wm^{-2} ($\sim 2\%$), based on deviations from monthly, meridional-mean LW IRF between 2001 and 2024. This uncertainty range is valid as long as Earth’s annual-mean climate state remains

within the seasonal and meridional variability sampled by ERA5LBL, the envelope 172
from which the regression coefficients were derived. Thus, the OLR-regression method 173
enables the LW IRF to be estimated directly from observed OLR, independent of 174
biases in atmospheric state variables and cloud radiative effects. 175

We apply this regression model to CERES EBAF v4.2 satellite observations of 176
OLR [57] to produce observationally-constrained estimates of LW IRF relative to 177
preindustrial levels. The OLR in the CERES EBAF product uses an objective con- 178
straint algorithm [58] that adjusts radiation fluxes for consistency with heat 179
storage in the Earth-atmosphere system; the remaining global-mean OLR uncer- 180
tainty of $\pm 2.5 \text{ Wm}^{-2}$ [57]. This yields annual time series of LW IRF for individual 181
WMGHGs over 2001–2024 (Fig.2). Over this 24-year period, the CERES-constrained 182
LW IRF for WMGHGs increases from **2.66 to 3.71** Wm^{-2} (2.66 to 3.66 Wm^{-2} if the 183
background climate is fixed at 2001). The 5-95% uncertainty range is **0.073** Wm^{-2} 184
($\sqrt{0.064^2 + 0.015^2 + 0.032^2}$, including 0.064 Wm^{-2} from the regression method, 0.015 185
 Wm^{-2} from clear-sky line-by-line calculations, and 0.032 Wm^{-2} from the uncer- 186
tainty in OLR; see Fig. S5 and S6 for a breakdown of the uncertainty related to each 187
WMGHG). 188

These observationally constrained estimates enable a direct comparison with LW 189
IRF derived independently from the cloud effects represented by ERA5LBL. Because 190
OLR in ERA5LBL is highly consistent with CERES EBAF v4.2, the CERES- 191
constrained values achieve excellent agreement with those computed from ERA5LBL. 192
This agreement reinforces confidence in our ERA5LBL calculations for LW IRF. 193

Implications to Earth System Modeling 194

In Earth system models (ESMs), effective radiative forcing (ERF) with fixed sea sur- 195
face temperature (F_{fsst}) is diagnosed as the TOA energy imbalance following rapid 196

197 atmospheric adjustments to greenhouse gas perturbations. These values are com-
 198 monly reported for $2\times\text{CO}_2$ and $4\times\text{CO}_2$ scenarios and serve as essential metrics for
 199 quantifying equilibrium climate sensitivity [7]. In IPCC AR6 [7], F_{fssst} was reported
 200 as $3.71 \pm 0.44 \text{ Wm}^{-2}$ for $2\times\text{CO}_2$ and $7.98 \pm 0.76 \text{ Wm}^{-2}$ for $4\times\text{CO}_2$, based on
 201 multi-model CMIP6 results [59].

202 In some ESMs, the IRF can be diagnosed using a second radiation call at every
 203 model timestep, for which absorber concentrations are perturbed while all other model
 204 fields remain unchanged. The fluxes from this diagnostic-only “double-call” do not
 205 feed back on the model physics and represent the instantaneous response to the spec-
 206 ified forcing agent. For WMGHG, such double-call perturbations are available only
 207 for $4\times\text{CO}_2$ in AMIP experiments [60], and only for a limited subset of six ESMs. As
 208 a result, $4\times\text{CO}_2$ is currently the only scenario for which IRF can be directly diag-
 209 nosed in these models. Because ESMs must rely on computationally efficient radiation
 210 parameterizations rather than spectrally resolved line-by-line radiative transfer, such
 211 double-call diagnostics provide a critical means of evaluating how well the internal
 212 radiation scheme of each model captures the intended forcing response.

213 Using the double-call diagnostics from the six ESMs and corresponding ERF values
 214 from AMIP- $4\times\text{CO}_2$ simulations (Table S6), we find that ERF ($8.10 \pm 1.12 \text{ Wm}^{-2}$)
 215 strongly correlates with LW IRF ($8.75 \pm 1.26 \text{ Wm}^{-2}$, $R^2 = 0.91$, Fig.3a), but shows a
 216 much weaker correlation with IRF_{TOA} ($R^2 = 0.48$, Table S6). Thus the LW IRF is a
 217 dominant cause of spread in ERF across these models.

218 To trace the source of inter-model spread in LW IRF, we conduct offline radia-
 219 tive transfer simulations using monthly-mean atmospheric conditions from each
 220 model, as described in the *Methods*. These simulations are performed under clear-
 221 sky conditions using a fast RTE-RRTMGP parameterization code [61], which isolates
 222 radiative responses from the model-specific representations of cloud processes. A
 223 known $+0.33 \text{ Wm}^{-2}$ bias in RTE-RRTMGP has been removed following benchmark

studies [14]. The resulting LW IRF values exhibit good agreement across models (black 224 points, x-axis of Fig.3b), with a spread of only 0.2 Wm^{-2} , much smaller than the 225 1.26 Wm^{-2} spread seen in double-call diagnostics. Importantly, this spread is not 226 driven by clouds: the variability under clear-sky double-call diagnostics (blue) is nearly 227 identical to that in all-sky (yellow), indicating that discrepancies arise mainly from 228 differences in radiation parameterizations. 229

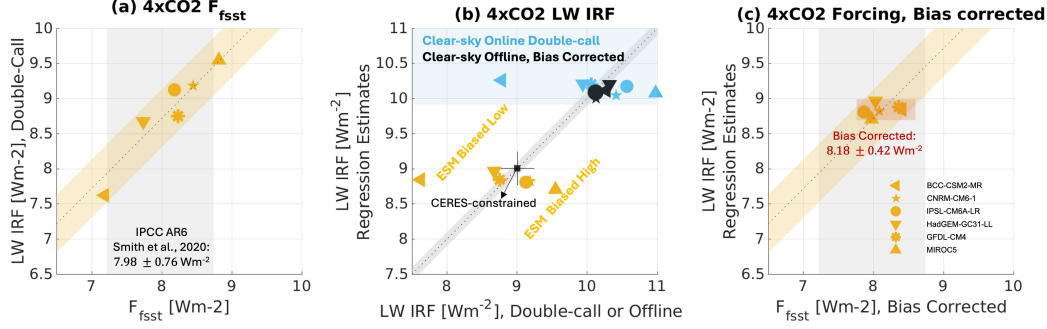


Fig. 3 Evaluation of $4\times\text{CO}_2$ radiative forcing in Earth System Models (ESMs). (a) Comparison of $4\times\text{CO}_2$ effective radiative forcing (F_{fsst} , x-axis) and longwave instantaneous radiative forcing at tropopause (LW IRF, y-axis) from online double radiation call diagnostics. (b) Comparison of LW IRF simulated by online double radiation call (x-axis; yellow: all-sky, blue: clear-sky), clear-sky offline RTE-RRTMGP with 0.33 Wm^{-2} bias correction applied (x-axis, black) [14], and estimates from the regression method applied to the simulated OLR (y-axis). CERES-constrained LW IRF from 8.95 to 9.06 Wm^{-2} over the 2001 to 2024 time period is shown as a black square. (c) Same as (a), but with model-specific bias corrections to LW IRF based on panel (b).

Although line-by-line calculations under all-sky conditions could, in principle, provide benchmark LW IRF estimates for each model, such calculations are computationally prohibitive and cannot realistically represent the model-specific cloud–radiation interaction processes. Instead, our regression framework offers a computationally efficient alternative, enabling accurate estimates of LW IRF across a broad range of atmospheric and cloud states as described by the OLR of these models.

Similar to previous sections, we derive the regression model for LW IRF $4\times\text{CO}_2$ based on clear-sky ERA5LBL calculations of OLR at pre-industrial CO_2 (CO_2PI),

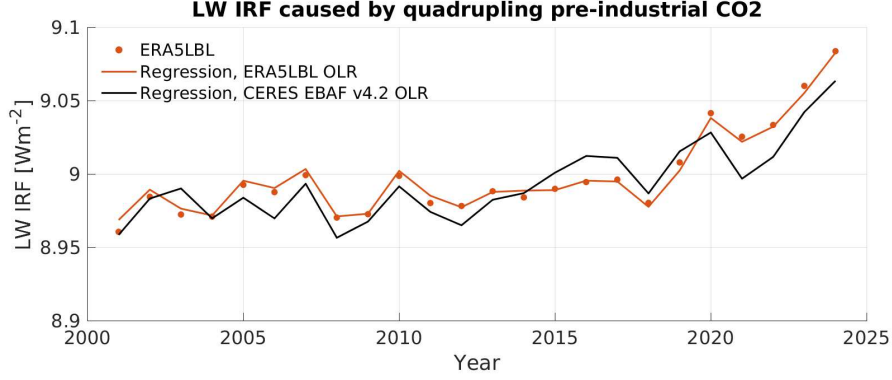


Fig. 4 LW IRF of $4\times\text{CO}_2$ (with respect to pre-industrial levels) at evolving background states from 2001 to 2024 from ERA5LBL (red scatters), the estimates from the regression applied to ERA5LBL OLR (red curve), and the estimates from the regression applied to CERES EBAF v4.2 OLR (black).

238 *Methods*) and calculations of LW IRF from quadrupled pre-industrial CO_2 ($\text{'CO}_24\text{xPI'}$,
 239 *Methods*) as:

$$F_{4\times\text{CO}_2} = 0.0509(R_{\text{CO}_2\text{PI}} - 63 \text{ Wm}^{-2}) - r, \quad (3)$$

240 where $R_{\text{CO}_2\text{PI}}$ represents the OLR if CO_2 were returned to 1850 levels while all other
 241 atmospheric variables remained at present-day conditions. This quantity is recon-
 242 structed as the sum of present-day OLR and the estimated $\text{LW IRF}_{\text{TOA}}$ from CO_2
 243 increase since 1850, for which the IRF_{TOA} is derived using the regression method
 244 (Table S5; evaluated in Figs. S4, S6). Our formulation uses pre-industrial levels,
 245 consistent with IPCC assessments. The global-mean bias evaluated at year 2010 is
 246 $r = 0.095 \text{ Wm}^{-2}$ (Table S2), and the 5–95% confidence interval is assessed to be
 247 $\pm 0.10 \text{ Wm}^{-2}$ using the 24-year monthly-mean, global-mean as described earlier. This
 248 regression-based uncertainty is smaller than the spread across five independent line-
 249 by-line models (0.19 Wm^{-2} , Table S1). In Fig.4, we apply this method to $R_{\text{CO}_2\text{PI}}$
 250 values derived from both ERA5LBL and CERES and have achieved excellent agree-
 251 ment. When using CERES-based global-mean OLR, the regression method estimates
 252 that LW IRF increases from 8.96 Wm^{-2} in 2001 to $9.06 \pm 0.24 \text{ Wm}^{-2}$ in 2024, driven
 253 by record-high OLR observed in the past two years due to surface warming [15] and

weakened cloud LW effects. The accuracy of these estimates under anomalous OLR 254
conditions demonstrates the robustness of the regression method in a changing climate. 255

Differences between the regression estimates and IRF from model double-call diag- 256
nostics provide a direct measure of radiation parameterization biases in ESMs. We 257
show that the regression technique explains the clear-sky LW IRF simulated by the 258
bias-corrected clear-sky offline calculations very well (Fig.3b, black, for six experiments 259
with double-call diagnostics; Fig.S7 for experiments without double-call). Extending 260
the regression to all-sky OLR from the double-call diagnostics of ESMs (Fig.3b, y- 261
axis), we find that the all-sky LW IRF of these models is $8.87 \pm 0.16 \text{ Wm}^{-2}$, close to 262
the CERES-constrained range of 8.95 to $9.06 \pm 0.24 \text{ Wm}^{-2}$. This reduced inter-model 263
spread reflects that ESMs agree well with observed global-mean OLR during the 264
historical period. Deviations from regression values are interpreted as model-specific 265
radiation biases. Assuming the reduced uncertainty in IRF can translate to an equiva- 266
lent reduction in ERF uncertainty, we derive a bias-corrected F_{fsst} of $8.18 \pm 0.42 \text{ Wm}^{-2}$ 267
(Fig.3c) for $4 \times \text{CO}_2$, significantly narrower than the original spread of 8.10 ± 1.12 268
 Wm^{-2} . This correction reduces inter-model spread in F_{fsst} by 60%. 269

The remaining uncertainty (yellow shading in Fig.3b) likely reflects differences in 270
shortwave forcing [13, 26, 49], adjustment processes [34, 59], and base-state conditions 271
[20]. As these processes are not fully independent of radiation, it is possible to achieve 272
a further reduction in inter-model spread if the same radiation parameterizations are 273
adopted. 274

Conclusion 275

This study establishes the first global benchmark for the LW IRF of WMGHGs under 276
realistic, all-sky conditions. Using a spectrally resolved, highly parallelized line-by-line 277
radiative transfer model, we quantify a post-1850 LW IRF of $3.71 \pm 0.07 \text{ Wm}^{-2}$ by 278
2024, with 38% of this increase occurring since 2001. Confidence in this benchmark 279

280 is reinforced by a regression method that integrates observational constraints from
281 satellite-observed OLR to account for uncertainties from cloud effects and evolving
282 climate states.

283 We further demonstrate that IRF values can be accurately derived from the OLR-
284 regression method applied to ESMs, offering a practical alternative to computationally
285 intensive line-by-line diagnostics for benchmarking climate projections. Applying this
286 framework reduces the uncertainty in CO₂-induced effective radiative forcing by over
287 half compared to the IPCC Sixth Assessment Report and provides a scalable pathway
288 for evaluating forcing across all major WMGHGs. A community-wide adoption of
289 this framework would strengthen the physical consistency of climate assessments and
290 enhance confidence in long-term climate projections.

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Methods

299

Line-by-line Calculations

300

The Geophysical Fluid Dynamics Laboratory’s GPU-compatible radiation code (GRT- 301
code) is a line-by-line radiative transfer code using HITRAN 2016 [10] and MT_CKD 302
2.5 [62]. It has been benchmarked in Pincus et al. (2020) [14], showing excellent 303
agreement with other line-by-line radiative transfer models including LBLRTM v12.8 304
[63] provided by Atmospheric and Environmental Research, RFM [64] calculation 305
conducted by Geophysical Fluid Dynamics Laboratory, 4AOP [65] provided by the 306
Laboratoire de Météorologie Dynamique, ARTS 2.3 [66] provided by the University of 307
Hamburg. We also note that although GRTcode does not include CO₂ line mixing, it 308
agrees well with models (such as RFM and LBLRTM) that include the mechanism, 309
which is consistent with the understanding that the line mixing induces negligible 310
(within 1%) impacts on the longwave radiative forcing [11]. Details of this comparison 311
are listed in Table S1. Inter-model discrepancies are estimated as the standard devi- 312
ations and are multiplied by 1.96 to infer 5-95% of uncertainty in IRF caused by gas 313
spectroscopy. 314

For clouds, the GRTcode adopts a high spectral resolution cloud optics parame- 315
terization developed by Feng et al., 2025 [56], which has been validated to achieve 316
reasonable radiative closure with hyperspectral resolution remote sensing instruments 317
in the mid-infrared [67, 68]. The cloud optics parameterization [56] is built upon Mie- 318
scattering calculations for liquid clouds and an optics library for severely roughened 319
solid columns developed by Yang et al. (2013) [69]. Based on grid-cell-mean cloud 320
fraction and cloud water content, subgrid cloud water content is stochastically gener- 321
ated at every wavenumber following Pincus et al., 2006 [30], and cloud optical depth 322
is assumed to be uniform within each 1 cm⁻¹ bin. The all-sky calculations of GRT- 323
code, with broadband cloud optics parameterizations, have shown good agreement in 324

Table S1 Clear-sky LW IRF [Wm^{-2}] from five line-by-line models [14] with respect to 100 profiles sampled from reanalysis products at year 2014. Standard deviations across these models are used to infer 5-95% uncertainty range in the line-by-line calculation of IRF. The OLR-regression method ('Regress') is applied using multi-model-mean clear-sky OLR and gas concentrations from Pincus et al., 2020 [14]. The ERA5LBL results for the global-mean of 2014 are listed based on an updated gas concentration datasets following CMIP7.

Scenario	4AOP	ARTS	GRTcode	RFM	LBLRTM	Mean/STD	Regress	ERA5LBL
OLR baseline	263.14	261.94	262.95	263.25	263.20	262.90/0.55	-	261.90
LW IRF at 200 hPa tropopause								
2014-1850 CO ₂	2.39	2.36	2.38	2.38	2.42	2.38/0.02	2.38	2.37
2014-1850 CH ₄	0.64	0.64	0.64	0.64	0.64	0.64/0.01	0.65	0.65
2014-1850 N ₂ O	0.21	0.21	0.21	0.21	0.21	0.21/0.01	0.21	0.22
2014-1850 CFC+HFC	0.42	0.42	0.42	0.41	0.38	0.41/0.02	0.43	0.43
4×CO ₂	10.26	10.15	10.20	10.23	10.40	10.25/0.10	10.22	10.17
LW IRF _{TOA}								
2014-1850 CO ₂	1.30	1.31	1.29	1.29	1.32	1.30/0.01	1.31	1.29
2014-1850 CH ₄	0.62	0.61	0.62	0.62	0.61	0.61/0.01	0.62	0.62
2014-1850 N ₂ O	0.21	0.21	0.20	0.21	0.21	0.21/0.01	0.21	0.21
2014-1850 CFC+HFC	0.57	0.56	0.56	0.55	0.50	0.55/0.03	0.56	0.56
4×CO ₂	5.40	5.45	5.32	5.37	5.51	5.41/0.08	5.38	5.33

long-term global-mean trend with space-borne hyper-spectral infrared sounders [40].

Using annual-mean, monthly-mean surface temperature, humidity, ozone concentration, and cloud conditions from the ERA5 [41] reanalysis dataset at 2.5° resolution, a set of global-scale line-by-line experiments perturbing the concentrations of well-mixed greenhouse gases (WMGHGs) was conducted for the period 2003–2022. These experiments include:

- Control: A control experiment using a time series of global-mean, annual-mean concentrations of WMGHGs [9].
- x PI: Identical to Control, except that an individual WMGHG species, x , is held fixed at its 1850 concentration (CO₂ : 284.297 parts per million by volume (ppmv), CH₄ : 0.7988 ppmv, N₂O : 0.2716 ppmv, 8.2×10^{-6} ppmv for CFC12eq and 2.02×10^{-5} ppmv HFC134a-eq).
- CO₂ 4xPI: Similar to CO₂ PI, except that CO₂ concentration is four times of 1850 concentration.

These experiments generate top-of-atmosphere outgoing longwave radiation (OLR, 339 Wm^{-2}). For the period from 2001 to 2022, the time series of WMGHG concentrations 340 is based on the Climate Model Intercomparison Project Phase 7 (CMIP7) greenhouse 341 input dataset [9]. For the years 2023 and 2024, global-mean concentrations of CO_2 , 342 CH_4 , and N_2O are taken from NOAA’s Global Monitoring Laboratory [42, 43], with 343 a bias correction applied to ensure continuity with Nicolls et al., 2025 [9]. Concen- 344 trations of CFC-12-eq and HFC-134a-eq gases after 2022 are held fixed at their 2022 345 levels. Radiative transfer calculations are conducted at 0.1 cm^{-1} resolution; we have 346 performed tests to demonstrate that increasing the resolution does not alter the broad- 347 band IRF. Table 1 lists the clear-sky LW IRF estimated from ERA5LBL calculations 348 for the year 2014, the difference between the GRTcode results submitted to Pincus et 349 al., 2020 [14] is minor and comes from samplings and differences in gas concentrations 350 that have been updated since CMIP6. 351

The Linear OLR-Regression Method 352

We find that the monthly-mean, gridded LW IRF of WMGHGs is largely a linear 353 function of OLR (Fig.S3), with linear coefficients closely matching those derived from 354 cloud-free (‘clear-sky’) simulations. Using the ERA5LBL simulation for the year 2010 355 and a first-order linear regression model (Equation 2), we derive coefficients (a and b) 356 from clear-sky simulations for both LW IRF and LW IRF at the top of the atmosphere 357 (IRF_{TOA} , Fig.S4), yielding zero global-mean bias (r). The global-mean bias under all- 358 sky conditions is then obtained by applying these coefficients to the all-sky OLR of the 359 same year. For LW IRF resulting from increases in WMGHG concentrations relative 360 to 1850 levels, the bias is within 0.02 Wm^{-2} ; for $4\times\text{CO}_2$, the bias is 0.09 Wm^{-2} . 361 These coefficients and biases are summarized in Table S2 (LW IRF) and Table S5 (LW 362 IRF_{TOA}) and are held fixed for each forcing scenario. 363

Table S2 Regression coefficients for instantaneous longwave radiative forcing of WMGHGs at tropopause. The forcing follows the form $F = a(R - b) - r$, where R is the outgoing longwave radiation (OLR). The coefficient a is a function of gas concentration (ppmv), where C_0 and C are unperturbed and perturbed concentrations, following logarithmic dependence for CO_2 [70], square-root dependence for CH_4 and N_2O [13], and linear dependence for CFCs and HFCs [22]. The term r is the global-mean bias in all-sky IRF (0 Wm^{-2} in clear-sky). These regression coefficients are derived from clear-sky monthly-mean, gridded ERA5LBL results at year 2010 and are evaluated using all-sky monthly-mean meridional-mean, providing the 5-95% uncertainty of the regression method.

Gas	a	$b \text{ (Wm}^{-2}\text{)}$	$r \text{ (Wm}^{-2}\text{)}$	5-95%
<i>Applied to OLR with present-day WMGHG concentrations</i>				
CO_2	$0.0358 \ln(C/C_0) + 0.0014 \ln(C/C_0)^2$	63.0	0.015	1.2%
CH_4	$0.0089(\sqrt{C} - \sqrt{C_0}) - 0.0008(\sqrt{C} - \sqrt{C_0})^2$	110.0	0.007	4.0%
N_2O	$0.027(\sqrt{C} - \sqrt{C_0})$	100.0	0.004	3.3%
CFC-12eq	$2.77(C - C_0)$	124.0	0.016	6.6%
HFC-134aeq	$1.46(C - C_0)$	124.0	0.002	8.1%
Sum			0.017	
<i>Applied to OLR with pre-industrial CO_2 and present-day other WMGHGs</i>				
$4 \times \text{CO}_2$	0.0509	63.0	0.095	1.0%

Table S3 All-sky LW IRF [Wm^{-2}] caused by increases in major WMGHGs from 1850 to 2014. Estimates from Etminan et al. (2016) [13] and IPCC AR6 [7] are converted from net stratospherically adjusted radiative forcing (SARF) using conversion ratios in Table S3 (second column). Estimates from Byrne and Goldblatt (2014) [12] are converted from net IRF using Table S3 (third column). CERES-constrained LW IRF based on the OLR-regression methods are also shown for comparison.

Source	Gas	Net SARF [Wm^{-2}]	LW IRF [Wm^{-2}]	Constrained LW IRF [Wm^{-2}]
Etminan et al. (2016) [13]	CO_2	1.806	2.059	2.097 ± 0.05
	CH_4	0.566	0.521	0.548 ± 0.02
	N_2O	0.175	0.173	0.187 ± 0.006
IPCC AR6 Ch.7 Supp. [7]	CFC12-eq	0.358	0.322	0.306 ± 0.039
	HFC134a-3q	0.040	0.036	0.038 ± 0.005
Byrne & Goldblatt (2014) [12]		Net IRF [Wm^{-2}]		
	CO_2	1.830	1.922	2.097 ± 0.05
	CH_4	0.525	0.515	0.548 ± 0.02
	N_2O	0.198	0.198	0.187 ± 0.006

364 The linear regression model explains more than 90% of the monthly gridded vari-
365 ability in LW IRF and LW IRF_{TOA}, except for LW IRF_{TOA} of CO_2 (Fig.S3-S4). This
366 exception arises because LW IRF_{TOA} is sensitive to stratospheric temperature [19–
367 21], which exhibit strong seasonal variations driven by the Brewer–Dobson circulation

Table S4 Conversion ratios from SARF or net IRF to LW IRF for major WMGHGs. References are listed in the left column for each gas.

Gas	LW IRF / Net SARF	LW IRF / Net IRF
CO ₂ [13, 49]	1.14	1.05
CH ₄ [13, 49]	0.92	0.98
N ₂ O [13, 14, 49]	0.99	1.00
CFC [4, 48]	0.90	–
HFC [4, 48]	0.90	–

[71]. On the other hand, the LW IRF used in this study, defined at the tropopause, is largely insensitive to stratospheric temperature and is therefore better captured by the regression model. This tropopause-based IRF is more closely aligned with the radiative forcing after stratospheric adjustment, which is primarily governed by radiative balance and provides a more robust indicator of surface climate change [6]. It has been adopted as the key metric for IRF since the first Intergovernmental Panel on Climate Change (IPCC) report [2, 4, 5]. Its importance is underscored by the strong correlation with the effective radiative forcing under fixed sea surface temperature (F_{fsst}), showing an R^2 of 0.91 (Fig.3a, Table S4). In contrast, the correlation between LW IRF_{TOA} and F_{fsst} is much weaker, with an R^2 of 0.48 (Table S4).

The linear regression model is then applied to OLR from ERA5LBL for 24 years, accurately reproducing both spatial and zonal-mean patterns in LW IRF for each WMGHG (Figs. S1-S2). For CO₂ -induced LW IRF_{TOA}, for which state-dependent analytical models have been developed [17, 19, 72], our linear model explains 97% of the annual-mean gridded variability, outperforming more complex models. For monthly-mean meridional mean LW IRF, the 5–95% uncertainty range of the OLR-regression method for each WMGHG (Figs. S5-S6) generally falls within 2% of the IRF, which is comparable with the spread of line-by-line calculations as reported by Pincus et al. (2020) [14]. The total uncertainty (σ) when using observed OLR is computed as:

$$\sigma = \sqrt{\sigma_r^2 + a^2\sigma_o^2 + \sigma_s^2} \quad (1)$$

where σ_r is the regression uncertainty, σ_o is the uncertainty from CERES EBAF v4.2 OLR, and σ_s is the line-by-line calculation uncertainty as summarized in Table S2 [14]. These are assumed to be independent and are combined in quadrature. Sources of uncertainty are indicated in Figs.2, S5-S6. While the annual mean and 4xCO₂ LW IRF have been presented in the main text, the LW IRF_{TOA} is evaluated to be 2.15±0.08 Wm⁻² for the 24-year average of WMGHG forcing relative to 1850 and 4.18±0.26 Wm⁻² for 4xCO₂.

Table S5 Similar to Table 1 but for LW IRF_{TOA}.

Gas	a	b (Wm ⁻²)	r (Wm ⁻²)
<i>Applied to OLR with present-day WMGHG concentrations</i>			
CO ₂	0.0359 ln(C/C_0)	155.0	0.007
CH ₄	0.0118($\sqrt{C} - \sqrt{C_0}$)ppmv ^{-1/2}	148.0	-0.010
N ₂ O	0.0366($\sqrt{C} - \sqrt{C_0}$)ppmv ^{-1/2}	148.0	-0.003
CFC-12eq	4.25($C - C_0$)ppmv ⁻¹	148.0	-0.001
HFC-134aeq	2.40($C - C_0$)ppmv ⁻¹	148.0	0.003
Sum			-0.014
<i>Applied to OLR with pre-industrial CO₂ and present-day other WMGHGs</i>			
4 × CO ₂	0.0493	155.0	0.108

RTE-RRTMGP Calculations with ESM inputs

RTE-RRTMGP [61] is a Fortran library for fast global-scale radiative transfer calculations. As described in Feng et al., (2023) [29], radiative transfer simulations at every model grid are conducted with inputs of monthly-mean surface temperature, temperature and humidity profiles from standard CMIP6 outputs at 19 pressure levels, climatological ozone profile, and present-day WMGHG concentrations except for

CO₂. Simulations are conducted with 1850 CO₂ concentrations and 4xCO₂ concentrations for 28 model runs conducted by 11 CMIP6 models for three experiments (piClim-control, piClim-4xCO₂, AMIP [60]), providing level-by-level downwelling and upwelling longwave radiative fluxes. These radiative fluxes are used to interpret the global-mean clear-sky OLR and LW IRF at each vertical level.

The radiative biases of RTE-RRTMGP can be evaluated using RFMIP [14] by comparing with line-by-line simulations for the same set of atmospheric profiles. The clear-sky LW IRF of 4xCO₂ simulated by RTE-RRTMGP is 10.536 Wm⁻², 0.33 Wm⁻² higher than that by GRTcode for the same set of atmospheric profiles in Pincus et al., 2020 [14]. This bias is removed in Figure 3(b) and Figure S7 when compared to IRF estimated by the OLR-regression method. While results for AMIP experiments are included in Figure 3 as black, other experiments are included in Figure S7 further confirming that the OLR-regression method is able to estimate the clear-sky LW IRF accurately within the proposed uncertainty range (0.1 Wm⁻², based on evaluations conducted in Figure S5).

Table S6 Radiation flux diagnostics [Wm⁻²] based on online double-radiation call of AMIP of six models, including BCC-CSM2-MR, CNRM-CM6-1, IPSL-CM6A-LR, MIROC5, GFDL-CM4, and HadGEM-GC31-LL. F_{fsst} is calculated as the net TOA radiation difference between AMIP and AMIP-4xCO₂ of the same model.

Variable	BCC	CNRM	IPSL	MIROC5	GFDL	HadGEM	Mean / Std
F_{fsst}	7.18	8.44	8.18	8.81	8.23	7.73	8.10 / 0.57
OLR	238.65	238.20	237.94	235.89	238.54	241.04	238.38 / 1.50
LW IRF _{TOA}	6.72	4.23	4.09	4.39	3.50	4.08	4.50 / 1.13
$F_{fsst} = \text{LW IRF} + \text{SW IRF} + \text{Adj}$							
LW IRF	7.62	9.19	9.13	9.55	8.75	8.67	8.82 / 0.67
SW IRF	-0.14	-0.69	-0.84	-0.56	-0.48	-0.39	-0.52 / 0.24
Adj	-0.30	-0.06	-0.11	-0.17	-0.04	-0.55	-0.20 / 0.19
LW IRF (Regression)	8.85	8.83	8.81	8.71	8.84	8.97	8.84 / 0.08
Radiation bias	1.23	-0.36	-0.31	-0.84	0.09	0.69	0.08 / 0.71
F_{fsst} (bias corrected)	8.41	8.08	7.87	7.98	8.32	8.42	8.18 / 0.21

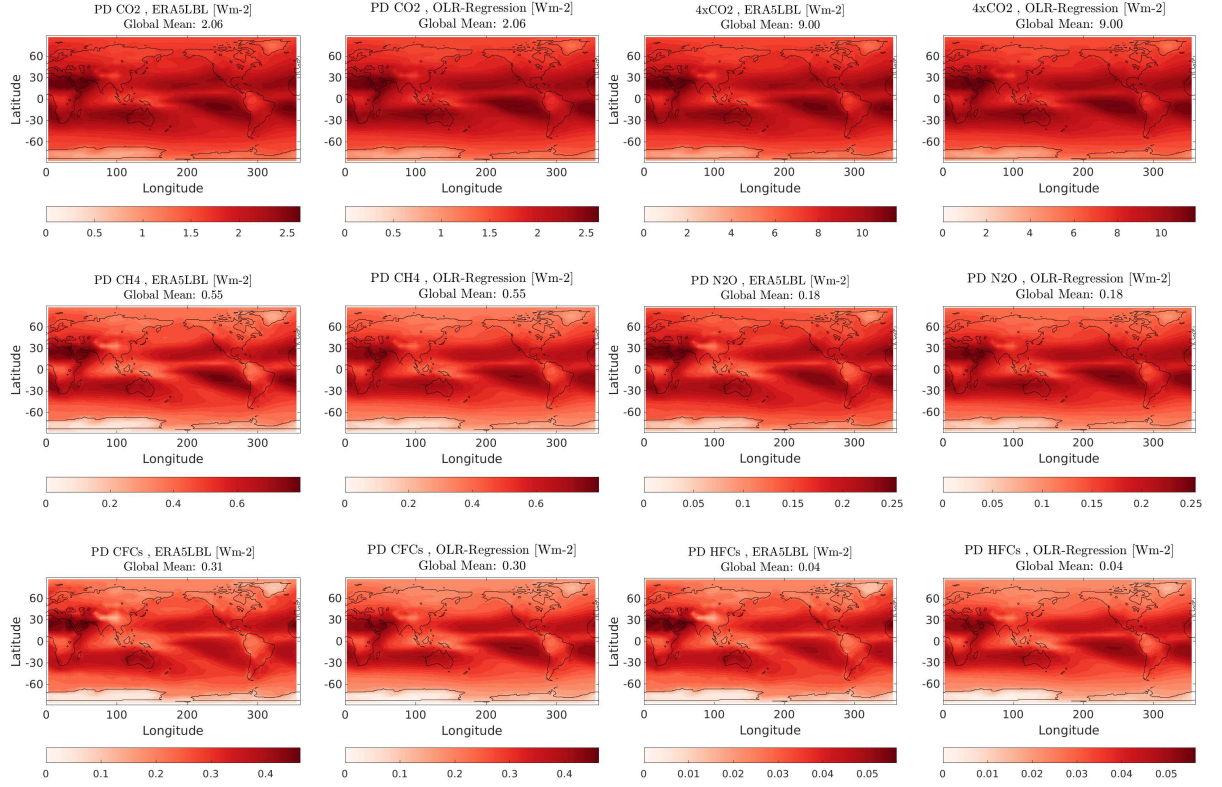


Fig. S1 Change in present-day (PD) (24-year mean, 2001 to 2024) LW IRF relative to the year 1850 caused CO₂, CH₄, N₂O, CFCs, HFCs, and 4xCO₂, from ERA5LBL and the regression estimates based on OLR simulated by ERA5LBL.

415 Declarations

416 • Data availability

417 Well-mixed greenhouse gas concentrations and radiative fluxes from ERA5LBL
 418 are available for all experiments listed in Supplementary Information at
 419 <https://doi.org/10.5281/zenodo.15734292>. CMIP6 and CMIP5 experiments are
 420 accessible via <https://aims2.llnl.gov/search> and summarized in *Methods*.

421 • Code availability

422 The GFDL line-by-line radiative transfer code is publicly accessible via
 423 <https://github.com/fengzydy/GRTCODE/tree/All-skyIRF>.

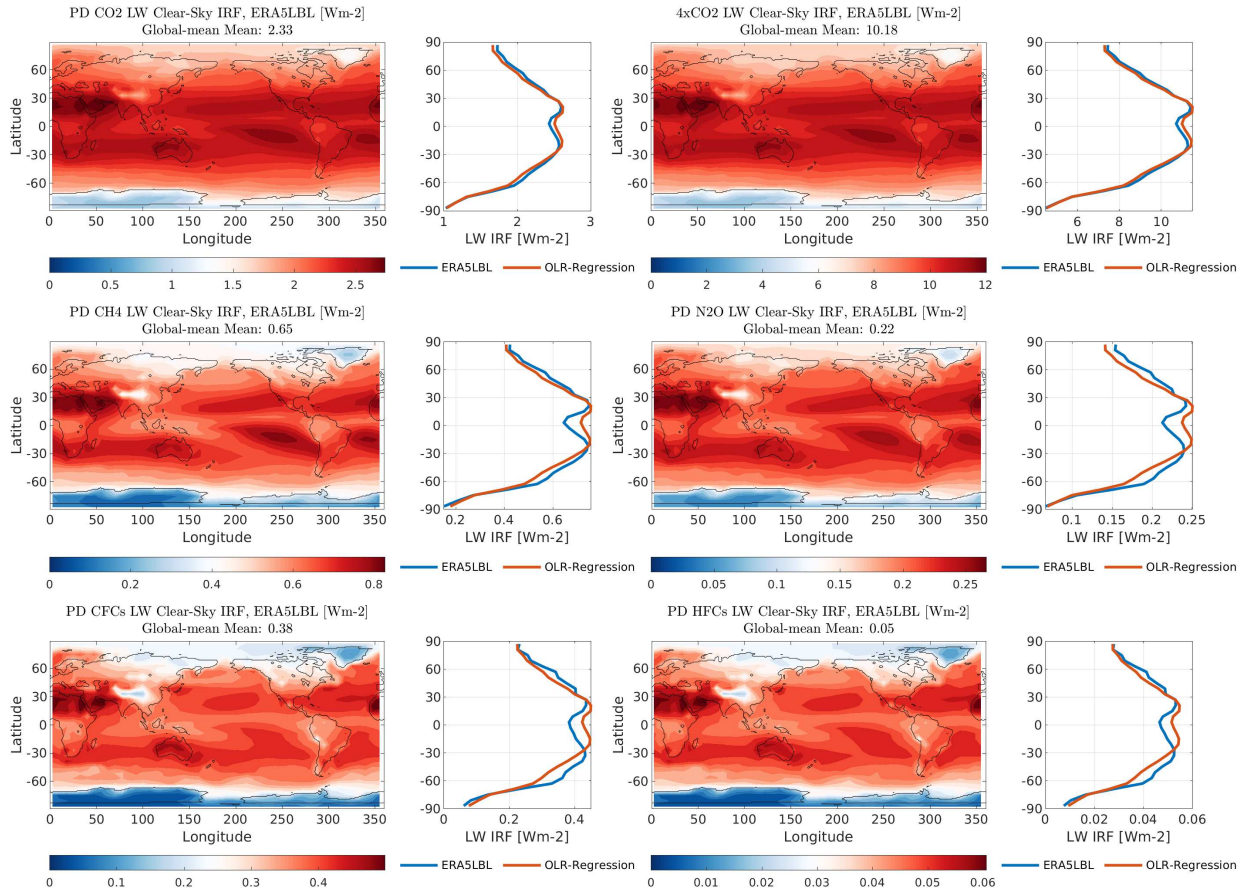


Fig. S2 The same as Figure S1 except for clear-sky IRF. Right panels of each figure compares the zonal-mean clear-sky obtained from the ERA5LBL (blue) and the OLR-regression method (red).

• Author contribution

424

J.F. and D.P. conceived the research project. J.F. conducted the line-by-line calcula- 425
 tions and designed the regression framework. R.M. and J.F. developed the GRTcode. 426
 J.F. and R.K. contributed to data processing and analyzes. All authors contributed 427
 to the manuscript writing. 428

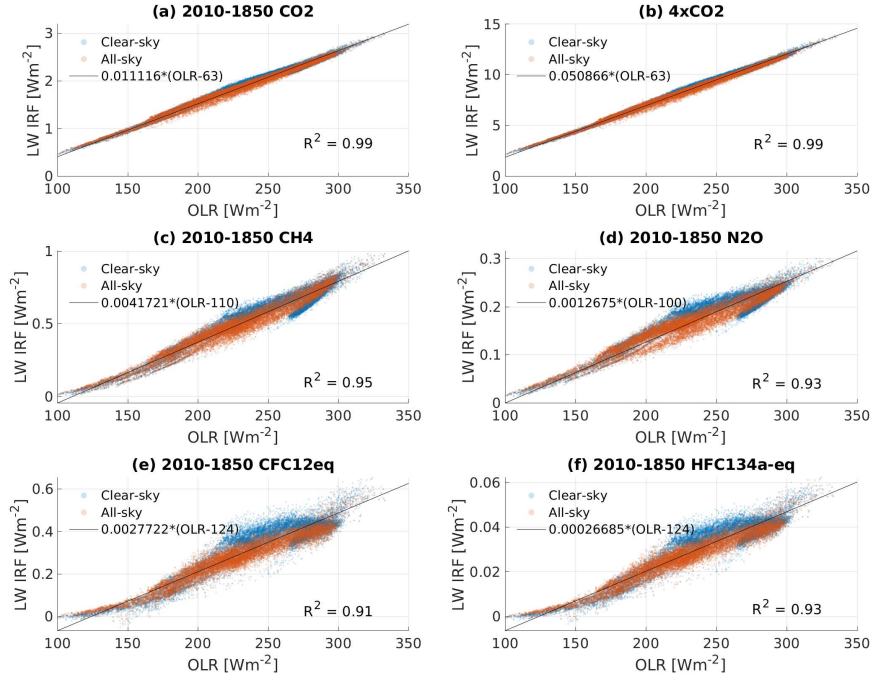


Fig. S3 Monthly gridded LW IRF as a function of OLR of the same location for (a) CO₂, (b) 4×CO₂, (c) CH₄, (d) N₂O, (e) CFC-12 equivalent, and (f) HFC-134a equivalent, each relative to preindustrial (1850) concentrations. Linear regression coefficients (Table S1) are derived from clear-sky ERA5LBL (blue), while biases and R^2 values are evaluated using all-sky ERA5LBL (red). Each scatter point represents a monthly-mean value for a grid cell from the ERA5LBL calculations at the year 2010.

References

- [1] Fels, S., Mahlman, J., Schwarzkopf, M. & Sinclair, R. Stratospheric sensitivity to perturbations in ozone and carbon dioxide- radiative and dynamical response. *Journal of the Atmospheric Sciences* **37**, 2265–2297 (1980).
- [2] Shine, K. P., Derwent, R., Wuebbles, D., Morcrette, J.-J. & Apling, A. in *Radiative forcing of climate* 41–68 (Cambridge University Press, 1990).
- [3] Griggs, D. J. & Noguera, M. Climate change 2001: the scientific basis. contribution of working group i to the third assessment report of the intergovernmental panel

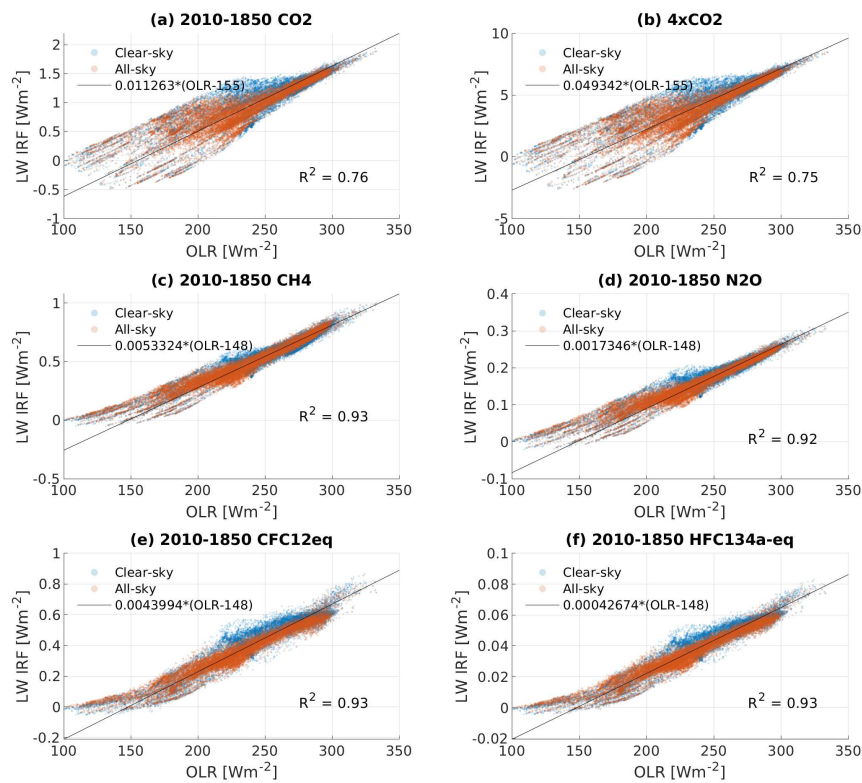


Fig. S4 Same as Fig.S3 but for LW IRF_{TOA}.

on climate change. *Weather* **57**, 267–269 (2002).

437

- [4] Collins, W. *et al.* Radiative forcing by well-mixed greenhouse gases: Estimates 438
from climate models in the intergovernmental panel on climate change (ipcc) 439
fourth assessment report (ar4). *Journal of Geophysical Research: Atmospheres* 440
111 (2006). 441

- [5] Myhre, G. *et al.* in *Anthropogenic and natural radiative forcing* (eds Stocker, 442
T. F. *et al.*) *Climate Change 2013: The Physical Science Basis. Contribution of* 443
Working Group I to the Fifth Assessment Report of the Intergovernmental Panel 444

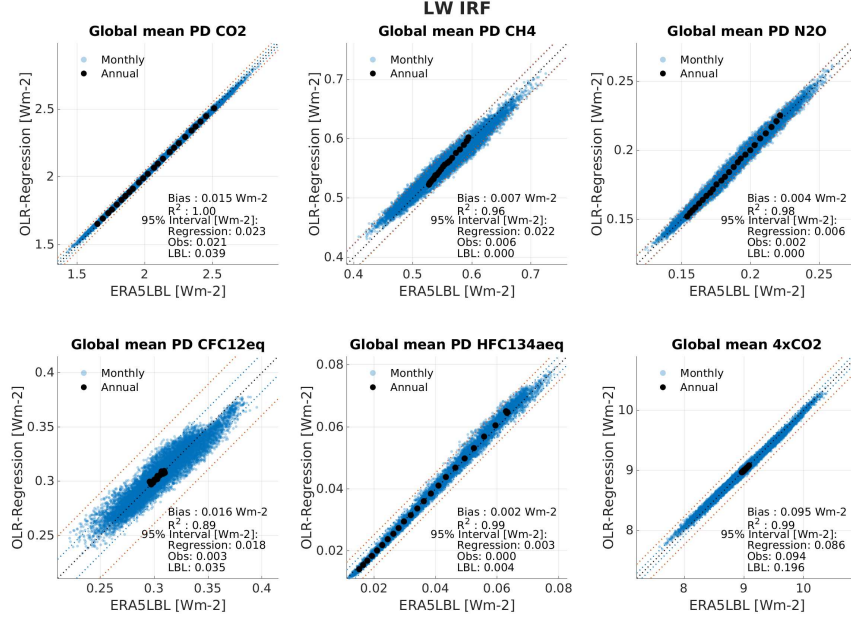


Fig. S5 Monthly meridional-mean (blue) and annual global-mean (black) LW IRF from 2001 to 2024, estimated using ERA5LBL (x-axis) and the OLR-regression method (y-axis) for individual WMGHG. Global-mean biases, R^2 , and uncertainties from the OLR-regression method, clear-sky line-by-line calculations, and observation of OLR, are marked in each panel.

- 445 on *Climate Change* 659–740 (Cambridge University Press, Cambridge, United
 446 Kingdom and New York, NY, USA, 2013). URL [https://www.ipcc.ch/report/](https://www.ipcc.ch/report/ar5/wg1/)
 447 [ar5/wg1/](https://www.ipcc.ch/report/ar5/wg1/).
- 448 [6] Ramaswamy, V. *et al.* Radiative forcing of climate: The historical evolution of
 449 the radiative forcing concept, the forcing agents and their quantification, and
 450 applications. *Meteorological Monographs* **59**, 14–1 (2019).
- 451 [7] Forster, P. *et al.* in *The earth's energy budget, climate feedbacks, and climate sen-*
 452 *sitivity. in climate change 2021: The physical science basis* (Cambridge University
 453 Press, 2021).

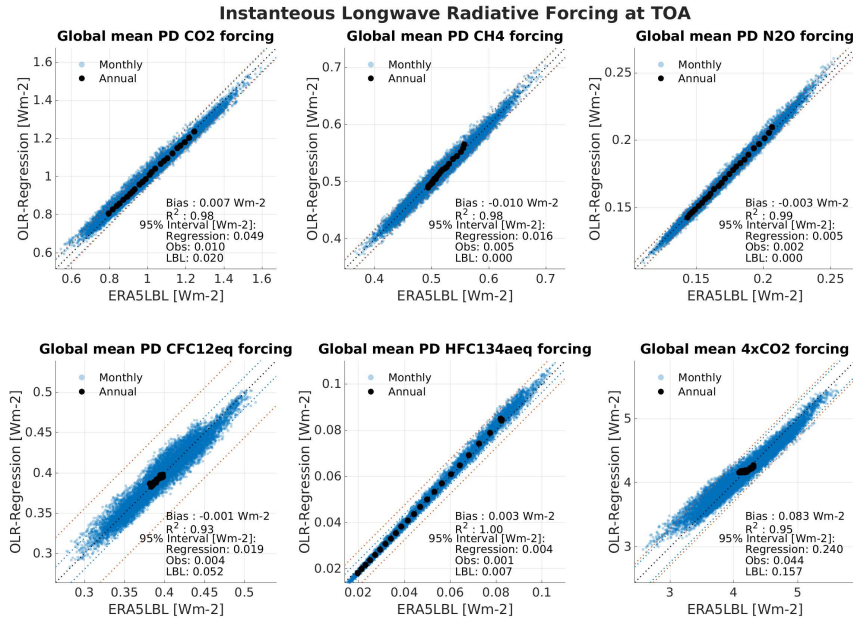


Fig. S6 Same as Fig.S5 for LW IRF at TOA.

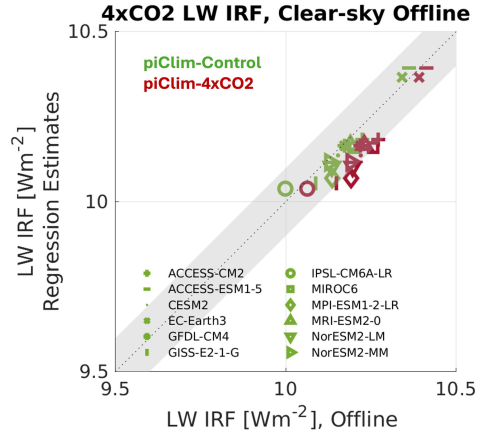


Fig. S7 The same as Figure 3(b) except that offline radiation calculations from RTE-RRTMGP driven by background states of piClim-Control and piClim-4xCO₂ experiments are compared to the OLR-regression method. The gray shading represents a 0.1 Wm⁻² uncertainty range of the OLR-regression method when applied to 4xCO₂ LW IRF.

- 454 [8] Meinshausen, M. *et al.* Historical greenhouse gas concentrations for climate
455 modelling (cmip6). *Geoscientific Model Development* **10**, 2057–2116 (2017).
- 456 [9] Nicholls, Z., Meinshausen, M., Pflüger, M. & Lewis, J. CMIP GHG Forcing
457 (v1.0.0) (2025). URL <https://doi.org/10.5281/zenodo.14892947>.
- 458 [10] Gordon, I. E. *et al.* The hitran2016 molecular spectroscopic database. *Journal*
459 *of quantitative spectroscopy and radiative transfer* **203**, 3–69 (2017).
- 460 [11] Mlynchak, M. G. *et al.* The spectroscopic foundation of radiative forcing of climate
461 by carbon dioxide. *Geophysical research letters* **43**, 5318–5325 (2016).
- 462 [12] Byrne, B. & Goldblatt, C. Radiative forcing at high concentrations of well-mixed
463 greenhouse gases. *Geophysical Research Letters* **41**, 152–160 (2014).
- 464 [13] Etminan, M., Myhre, G., Highwood, E. J. & Shine, K. P. Radiative forcing of
465 carbon dioxide, methane, and nitrous oxide: A significant revision of the methane
466 radiative forcing. *Geophysical Research Letters* **43**, 12–614 (2016).
- 467 [14] Pincus, R. *et al.* Benchmark calculations of radiative forcing by greenhouse gases.
468 *Journal of Geophysical Research: Atmospheres* **125**, e2020JD033483 (2020).
- 469 [15] Forster, P. M. *et al.* Indicators of global climate change 2024: annual update of
470 key indicators of the state of the climate system and human influence. *Earth*
471 *System Science Data* **17**, 2641–2680 (2025).
- 472 [16] Andrews, T. & Forster, P. M. Co2 forcing induces semi-direct effects with con-
473 sequences for climate feedback interpretations. *Geophysical Research Letters* **35**
474 (2008).
- 475 [17] Zhang, M. & Huang, Y. Radiative forcing of quadrupling co2. *Journal of Climate*
476 **27**, 2496–2508 (2014).

- [18] Huang, Y., Tan, X. & Xia, Y. Inhomogeneous radiative forcing of homogeneous 477
greenhouse gases. *Journal of Geophysical Research: Atmospheres* **121**, 2780–2789 478
(2016). 479
- [19] Jeevanjee, N., Seeley, J. T., Paynter, D. & Fueglistaler, S. An analytical model for 480
spatially varying clear-sky co₂ forcing. *Journal of Climate* **34**, 9463–9480 (2021). 481
- [20] He, H., Kramer, R. J., Soden, B. J. & Jeevanjee, N. State dependence of co₂ 482
forcing and its implications for climate sensitivity. *Science* **382**, 1051–1056 (2023). 483
- [21] Chen, Y.-T., Merlis, T. M. & Huang, Y. The cause of negative co₂ forcing at 484
the top-of-atmosphere: The role of stratospheric versus tropospheric temperature 485
inversions. *Geophysical Research Letters* **51**, e2023GL106433 (2024). 486
- [22] Forster, P. M. d. F. *et al.* Resolution of the uncertainties in the radiative forcing 487
of hfc-134a. *Journal of Quantitative Spectroscopy and Radiative Transfer* **93**, 488
447–460 (2005). 489
- [23] Myhre, G., Stordal, F., Gausemel, I., Nielsen, C. J. & Mahieu, E. Line-by-line 490
calculations of thermal infrared radiation representative for global condition: Cfc- 491
12 as an example. *Journal of Quantitative Spectroscopy and Radiative Transfer* 492
97, 317–331 (2006). 493
- [24] Hodnebrog, Ø. *et al.* Global warming potentials and radiative efficiencies of halo- 494
carbons and related compounds: A comprehensive review. *Reviews of Geophysics* 495
51, 300–378 (2013). 496
- [25] Maycock, A. C., Smith, C. J., Rap, A. & Rutherford, O. On the structure 497
of instantaneous radiative forcing kernels for greenhouse gases. *Journal of the* 498
Atmospheric Sciences **78**, 949–965 (2021). 499

- 500 [26] Byrom, R. E. & Shine, K. P. Methane’s solar radiative forcing. *Geophysical*
501 *Research Letters* **49**, e2022GL098270 (2022).
- 502 [27] Czarnecki, P., Polvani, L. & Pincus, R. Analytical models of instantaneous
503 radiative forcing across opacity regimes. *Authorea Preprints* (2025).
- 504 [28] Koll, D. D. & Cronin, T. W. Earth’s outgoing longwave radiation linear due to
505 h₂o greenhouse effect. *Proceedings of the National Academy of Sciences* **115**,
506 10293–10298 (2018).
- 507 [29] Feng, J., Paynter, D., Wang, C. & Menzel, R. How atmospheric humidity
508 drives the outgoing longwave radiation–surface temperature relationship and
509 inter-model spread. *Environmental Research Letters* **18**, 104033 (2023).
- 510 [30] Pincus, R., Hemler, R. & Klein, S. A. Using stochastically generated subcolumns
511 to represent cloud structure in a large-scale model. *Monthly weather review* **134**,
512 3644–3656 (2006).
- 513 [31] Davis, S. M. *et al.* Assessment of upper tropospheric and stratospheric water
514 vapor and ozone in reanalyses as part of s-rip. *Atmospheric chemistry and physics*
515 **17**, 12743–12778 (2017).
- 516 [32] Wright, J. S. *et al.* Differences in tropical high clouds among reanalyses: ori-
517 gins and radiative impacts. *Atmospheric Chemistry and Physics* **20**, 8989–9030
518 (2020).
- 519 [33] Huang, H. & Huang, Y. Diagnosing the radiation biases in global climate models
520 using radiative kernels. *Geophysical Research Letters* **50**, e2023GL103723 (2023).
- 521 [34] Smith, C. J. *et al.* Effective radiative forcing and adjustments in cmip6 models.
522 *Atmospheric Chemistry and Physics* **20**, 9591–9618 (2020).

- [35] Cess, R. *et al.* Uncertainties in carbon dioxide radiative forcing in atmospheric 523
general circulation models. *Science* **262**, 1252–1255 (1993). 524
- [36] Strow, L. L. & DeSouza-Machado, S. Establishment of air's climate-level radio- 525
metric stability using radiance anomaly retrievals of minor gases and sea surface 526
temperature. *Atmospheric Measurement Techniques* **13**, 4619–4644 (2020). 527
- [37] Teixeira, J., Wilson, R. C. & Thrastarson, H. T. Direct observational evidence 528
from space of the effect of co₂ increase on longwave spectral radiances: the 529
unique role of high-spectral-resolution measurements. *Atmospheric Chemistry* 530
and Physics **24**, 6375–6383 (2024). 531
- [38] Chen, X., Huang, X. & Strow, L. L. Near-global cfc-11 trends as observed by 532
atmospheric infrared sounder from 2003 to 2018. *Journal of Geophysical Research:* 533
Atmospheres **125**, e2020JD033051 (2020). 534
- [39] Kramer, R. J. *et al.* Observational evidence of increasing global radiative forcing. 535
Geophysical Research Letters **48**, e2020GL091585 (2021). 536
- [40] Raghuraman, S. P., Paynter, D., Ramaswamy, V., Menzel, R. & Huang, X. Green- 537
house gas forcing and climate feedback signatures identified in hyperspectral 538
infrared satellite observations. *Geophysical Research Letters* **50**, e2023GL103947 539
(2023). 540
- [41] Hersbach, H. *et al.* The era5 global reanalysis. *Quarterly Journal of the Royal* 541
Meteorological Society **146**, 1999–2049 (2020). 542
- [42] Lan, X., Tans, P., Thoning, K. *et al.* Trends in globally-averaged co₂ determined 543
from noaa global monitoring laboratory measurements. *NOAA GML* (2023). 544

- 545 [43] Lan, K. T., X. & Dlugokencky, E. Trends in globally-averaged ch4, n2o, and sf6
546 determined from noaa global monitoring laboratory measurements.version 2025-
547 05. *NOAA GML* (2025).
- 548 [44] Ramanathan, V., Lian, M. & Cess, R. Increased atmospheric co2: Zonal and
549 seasonal estimates of the effect on the radiation energy balance and surface
550 temperature. *Journal of Geophysical Research: Oceans* **84**, 4949–4958 (1979).
- 551 [45] WMO. Report of the meeting of experts on potential climatic effects of ozone
552 and other minor trace gases. Tech. Rep. (1982).
- 553 [46] Soden, B. J. *et al.* Quantifying climate feedbacks using radiative kernels. *Journal*
554 *of Climate* **21**, 3504–3520 (2008).
- 555 [47] Vial, J., Dufresne, J.-L. & Bony, S. On the interpretation of inter-model spread
556 in cmip5 climate sensitivity estimates. *Climate Dynamics* **41**, 3339–3362 (2013).
- 557 [48] Shine, K. P. & Myhre, G. The spectral nature of stratospheric temperature adjust-
558 ment and its application to halocarbon radiative forcing. *Journal of Advances in*
559 *Modeling Earth Systems* **12**, e2019MS001951 (2020).
- 560 [49] Shine, K. P., Byrom, R. E. & Checa-Garcia, R. Separating the shortwave and
561 longwave components of greenhouse gas radiative forcing. *Atmospheric Science*
562 *Letters* **23**, e1116 (2022).
- 563 [50] Oreopoulos, L. & Cahalan, R. F. Cloud inhomogeneity from modis. *Journal of*
564 *climate* **18**, 5110–5124 (2005).
- 565 [51] Barker, H. W., Stephens, G. L. & Fu, Q. The sensitivity of domain-averaged
566 solar fluxes to assumptions about cloud geometry. *Quarterly Journal of the Royal*
567 *Meteorological Society* **125**, 2127–2152 (1999).

- [52] Hogan, R. J. & Illingworth, A. J. Deriving cloud overlap statistics from radar. 568
Quarterly journal of the royal meteorological society **126**, 2903–2909 (2000). 569
- [53] Oreopoulos, L., Cho, N. & Lee, D. Revisiting cloud overlap with a merged dataset 570
of liquid and ice cloud extinction from cloudsat and calipso. *Frontiers in Remote* 571
Sensing **3**, 1076471 (2022). 572
- [54] Mie, G. Beiträge zur Optik trüber Medien, speziell kolloidaler Metallösungen. 573
Annalen der Physik **330**, 377–445 (1908). 574
- [55] Mitchell, D. L. Effective diameter in radiation transfer: General definition, appli- 575
cations, and limitations. *Journal of the Atmospheric Sciences* **59**, 2330–2346 576
(2002). 577
- [56] Feng, J., Menzel, R. & Paynter, D. Toward transparency and consistency: An 578
open-source optics parameterization for clouds and precipitation. *Journal of* 579
Advances in Modeling Earth Systems **17**, e2024MS004478 (2025). 580
- [57] Loeb, N. G. *et al.* Clouds and the earth’s radiant energy system (ceres) energy bal- 581
anced and filled (ebaf) top-of-atmosphere (toa) edition-4.0 data product. *Journal* 582
of Climate **31**, 895–918 (2018). 583
- [58] Loeb, N. G. *et al.* Toward optimal closure of the earth’s top-of-atmosphere 584
radiation budget. *Journal of Climate* **22**, 748–766 (2009). 585
- [59] Smith, C. *et al.* Understanding rapid adjustments to diverse forcing agents. 586
Geophysical Research Letters **45**, 12–023 (2018). 587
- [60] Eyring, V. *et al.* Overview of the coupled model intercomparison project phase 6 588
(cmip6) experimental design and organization. *Geoscientific Model Development* 589
9, 1937–1958 (2016). 590

- 591 [61] Pincus, R., Mlawer, E. J. & Delamere, J. S. Balancing accuracy, efficiency, and
592 flexibility in radiation calculations for dynamical models. *Journal of Advances in*
593 *Modeling Earth Systems* **11**, 3074–3089 (2019).
- 594 [62] Mlawer, E. J. *et al.* Development and recent evaluation of the mt_ckd model
595 of continuum absorption. *Philosophical Transactions of the Royal Society A:*
596 *Mathematical, Physical and Engineering Sciences* **370**, 2520–2556 (2012).
- 597 [63] Clough, S. A. *et al.* Atmospheric radiative transfer modeling: A summary of
598 the aer codes. *Journal of Quantitative Spectroscopy and Radiative Transfer* **91**,
599 233–244 (2005).
- 600 [64] Dudhia, A. The reference forward model (rfm). *Journal of Quantitative*
601 *Spectroscopy and Radiative Transfer* **186**, 243–253 (2017).
- 602 [65] Cheruy, F., Scott, N., Armante, R., Tournier, B. & Chedin, A. Contribution to the
603 development of radiative transfer models for high spectral resolution observations
604 in the infrared. *Journal of Quantitative Spectroscopy and Radiative Transfer* **53**,
605 597–611 (1995).
- 606 [66] Buehler, S. A. *et al.* Arts, the atmospheric radiative transfer simulator—version
607 2.2. *Geoscientific Model Development* **11**, 1537–1556 (2018).
- 608 [67] Feng, J. & Huang, Y. Impacts of tropical cyclones on the thermodynamic condi-
609 tions in the tropical tropopause layer observed by a-train satellites. *Atmospheric*
610 *Chemistry and Physics* **21**, 15493–15518 (2021).
- 611 [68] Liu, L., Huang, Y. & Gyakum, J. R. Clouds reduce downwelling longwave
612 radiation over land in a warming climate. *Nature* 1–7 (2025).

- [69] Yang, P. *et al.* Spectrally consistent scattering, absorption, and polarization 613
properties of atmospheric ice crystals at wavelengths from 0.2 to 100 μ m. *Journal* 614
of the atmospheric sciences **70**, 330–347 (2013). 615
- [70] Huang, Y. & Bani Shahabadi, M. Why logarithmic? a note on the depen- 616
dence of radiative forcing on gas concentration. *Journal of Geophysical Research:* 617
Atmospheres **119**, 13–683 (2014). 618
- [71] Young, P. J., Thompson, D. W., Rosenlof, K. H., Solomon, S. & Lamarque, J.-F. 619
The seasonal cycle and interannual variability in stratospheric temperatures and 620
links to the brewer–dobson circulation: An analysis of msu and ssu data. *Journal* 621
of Climate **24**, 6243–6258 (2011). 622
- [72] Chen, Y.-T., Huang, Y. & Merlis, T. M. The global patterns of instantaneous co 623
2 forcing at the top of the atmosphere and the surface. *Journal of Climate* **36**, 624
6331–6347 (2023). 625